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Proper Accounting of Mass Transfer Resistances in Forward Osmosis: Improving the Accuracy of Model Predictions of Structural Parameter

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Abstract

This work demonstrates a more accurate method for calculating structural parameter (S) of asymmetric osmotic membranes using experimental data and a theoretical flux model which encapsulates all significant boundary layer phenomena. External boundary layer effects on the porous side of the membrane have been neglected in many current models. In these models, external concentration polarization (ECP) effects get combined with the internal concentration polarization (ICP), resulting in inflated S values. In this study, we proposed a new flux model in which ECP effects are accounted for so that S can be more accurately measured. This model considered the in-series resistances for solute transport based on intrinsic properties of the membrane as well as boundary layers at membrane surfaces and within the support layer. The results indicate that ICP is less severe than previously predicted and that cross-flow velocity, temperature and concentration of the draw and the feed solutions impact both external and internal concentration polarization. Our calculations also surprisingly show that changes in cross-flow velocity impact internal concentration polarization due to induced mixing within the support layer. Also, we suggest that it is critical to consider the “residence time” of solutes in the vicinity of the selective layer in determining the membrane selectivity.
1. Introduction

Forward Osmosis (FO) and pressure retarded osmosis (PRO) have recently been revitalized as a sustainable and versatile membrane-based separation technology platform for water and power production, respectively. These processes are driven by the osmotic pressure difference between two solutions separated by a semi-permeable membrane. There are a number of configurations, uses and system integration methods proposed in the literature [1-3]. Some of these processes include treating impaired water and landfill leachate [4-7], recovery of high-value dissolved solids in food processing and pharmaceutical industries [8, 9], desalination [1, 10-12], or salinity-gradient power generation [3, 13-21]. FO has been widely studied in recent years, mainly focusing on development of improved membrane structures [22-32] and draw solutes [33-37] for efficient operation.

To design membranes for various FO processes, it is important to understand critical structure-performance relationships, especially with respect to mass transfer. This has led to a number of theoretical transport models to predict osmotic water flux coupled with reverse salt flux across FO membranes [21, 23, 38-50]. Many of these models have been used to predict the severity of concentration polarization (CP), a mass transfer phenomenon that reduces osmotic driving force. Specifically, internal concentration polarization (ICP), which mimics a stagnant diffusion limited boundary layer occurring within an asymmetric membrane support structure, has largely been recognized as the major impediment to membrane flux performance [38-40, 51]. Lee et. al. derived one of the first mathematical models describing transport during osmosis across asymmetric membranes and was among the first to describe ICP [39]. Still emphasizing the importance of ICP, McCutcheon and coworkers [38] developed a flux model in which the impact of external concentration polarization at the selective layer surface (ECPs) was combined. Phillip and coworker presented a model describing the reverse solute permeation across an asymmetric membrane in FO [43]. Recently, Yip et.al. [21] and Tiraferri et.al. [52]
developed models in which ICP, ECP\textsuperscript{p} and reverse salt permeation were all taken into account to predict flux performance in FO. In these models, although the osmotic flux behaviors appeared to be predicted accurately, the effect of ECP occurring on the porous side of the membrane structure (ECP\textsuperscript{p}) has commonly been ignored. This assumption was valid for early membranes that exhibited low flux or having large structural parameters relative to the external boundary layer thickness. However, with the advent of high performance membranes operated at higher flux or at low crossflow velocities, ECP\textsuperscript{p} becomes more substantial and the assumption is no longer valid \cite{50,62}.

Neglecting the ECP will result in an inaccurate accounting of the severity of ICP. If neglected, existing models simply lump ECP\textsuperscript{p} into the ICP term. This results in an overprediction of structural parameter since the ECP effects are misallocated to structural resistances in the support side of the membrane. Limited attention on ECP\textsuperscript{p} could misrepresent the severity of ICP and misdirect research priorities to support layer development when in fact ECP\textsuperscript{p} could become a more critical problem as better membranes become available.

A recent computational fluid dynamics simulation on forward osmosis systems using asymmetric membranes supported this hypothesis in revealing the non-negligible ECP\textsuperscript{p} even when high cross-flow velocity and slip velocity at the porous surface were accounted \cite{50}. If we are to continue using experimental data to characterize structural parameter, a comprehensive mass transfer model is needed to account for this ECP\textsuperscript{p} so that proper attention is paid to both membrane design and external fluid dynamics conditions. Improved understanding of this phenomenon is essential for better overall system modeling as pilot- and full-scale systems are built and cost models established.

This study demonstrates an updated theoretical model to predict osmotic flux performance of asymmetric and thin-film composite membranes in FO. ECP\textsuperscript{p} was incorporated
in such a flux model to advance understanding of its impacts to osmotic flux. The related mass transfer resistances, e.g. ICP, ECP\(^p\), ECP\(^s\), solute resistivity and structural parameter, were effectively de-convoluted from one another. This led to a new governing equation defining solute resistivity and structural parameter. Furthermore, new definitions of membrane reflection coefficient and total resistance to solute transport emerge from this work. “Residence time” of solutes at membrane surfaces is, for the first time, discussed as it relates to membrane selectivity. This model is supported by experimental data from two commercial forward osmosis membranes.

2. Theory and Model Development

Figure 1 shows a schematic of the solute concentration gradient across a thin-film composite membrane operating in FO mode (porous support layer faces a high-salinity draw solution) [62]. The effective driving force for osmosis exists only at the interfaces of the selective layer. From the bulk draw solution, solutes must diffuse through the porous support to the interface between the support and selective layers. Once there, a small amount of solute diffuses across the selective layer to the lower-saline feed [43]. Below, a new model is derived which accounts for ICP, ECP\(^s\) and ECP\(^p\) coupled with the reverse solute flux.

2.1. Updated Model: Incorporating ECP on Both Sides of the Membrane

2.1.1. For the FO Mode

Osmotic water flux across the selective layer, \(J_w\), is defined as:

\[
J_w = A(\sigma \Delta \pi_b - \Delta P) = A(\Delta \pi_m - \Delta P) = A\left[\left(\pi_{D,m} - \pi_{F,m}\right) - \Delta P\right]
\]

(1)

where, \(A\) is the intrinsic water permeability coefficient, or permeance, of the membrane, \(\sigma\) is the reflection coefficient of the membrane, \(\Delta \pi_m\) is the effective osmotic pressure difference across
the selective layer, $\Delta P$ is the transmembrane hydrostatic pressure. $\Delta P$ is typically zero or close to zero in many FO applications.

Reverse salt flux, $J_{s}^{a}$, across the selective layer:

$$J_{s}^{a} = B \Delta C_{m} = B(C_{D,m} - C_{F,m})$$

(2)

where, $B$ is the intrinsic solute permeability coefficient of the membrane, $C_{D,m}$ and $C_{F,m}$ are the solute concentrations at the selective interface in the draw and feed, respectively.

As water permeates (i.e. osmotic water flux, $J_{w}$) from the feed to the draw, it dilutes the draw solution within the porous support and draw solution boundary layer causing both dilutive ICP and ECP$^p$. At the same time, solutes diffuse from the bulk draw, through the porous support, to the interface of this support and the selective layer due to a concentration gradient from dilution of the draw solution. This transport is described by the commonly used convection-diffusion equation, consisting of a diffusive term caused by the solute concentration gradient and a negative convective term caused by the mass flow of water across the membrane in the opposite direction [39].

$$J_{s}^{a} = D \frac{dC(x)}{dx} - J_{w}C(x)$$

(3)

where, $D$ is the solute diffusion coefficient. $D$ is then referred to $D_{s}$ or $D_{D}$ when accounting for solute diffusivity inside the support layer or in the draw solution, respectively. In the porous support structure, $D_{s}$ is the bulk diffusion coefficient corrected with the porosity, $\varepsilon$, and tortuosity, $\tau$, of the support, or $D_{s} = \frac{D_{D}\varepsilon}{\tau}$.

At steady state, the solute fluxes across the selective layer and the porous support are equal:

$$J_{s}^{a} = J_{s}^{d}$$

(4)
or

\[ B(C_{D,m} - C_{F,m}) = D \frac{dC(x)}{dx} - J_w C(x) \]  

(5)

To solve the differential equation, we consider the following boundary conditions:

\[
\begin{align*}
    x &= 0, \quad C(x) = C_{D,i} \\
    x &= -t_s, \quad C(x) = C_{D,m}, \quad D = D_i = \frac{D_p \varepsilon}{\tau} \\
    x &= \delta_D, \quad C(x) = C_{D,b}, \quad D = D_D
\end{align*}
\]

Boundary conditions:

We also define

\[
    t_s = \frac{S \left( \frac{\varepsilon}{\tau} \right)}{D_D \left( \frac{\varepsilon}{\tau} \right)} = \frac{S}{D_D} \quad \text{and} \quad \delta_D = \frac{1}{k_D}
\]

where, \( S = t_s \tau / \varepsilon \) is the structural parameter of the membrane support, \( \delta_D \) and \( k_D \) are the thickness of the boundary layer and the mass transfer coefficient in the draw side, respectively.

Solving eqn. (5) for the boundary conditions, the draw solution concentration is defined by:

\[
    C_{D,m} = C_{D,b} \exp \left[ -J_w \left( \frac{1}{k_D} + \frac{S}{D_D} \right) \right] + B \Delta C_m \left\{ \exp \left[ -J_w \left( \frac{1}{k_D} + \frac{S}{D_D} \right) \right] - 1 \right\} \]  

(6)

In this equation, the draw solute concentration at the support – selective layer interface, \( C_{D,m} \), is the sum of two components on the right expression. The first component is the concentration of the bulk draw solution, \( C_{D,b} \), accounted for dilutive ECP\(^p\) and ICP effects which are described by the exponential terms of \( -\frac{J_w}{k_D} \) and \( -\frac{J_w S}{D_D} \). The second contribution to the increase of the solute concentration at the membrane surface is from the reverse solute permeation.

On the feed side, the convective water flux from the feed to the draw carries solutes from the bulk feed to the selective layer where they theoretically are rejected and accumulate. This is a concentrative external concentration polarization, with a boundary layer thickness of \( \delta_F = D_F k_F \), where \( D_F \) is solute diffusion coefficient and \( k_F \) is the mass transfer coefficient in the feed. The salt flux within this cECP\(^s\) layer is, also, a combination of the diffusive and convective terms:
\[ J_{i}^{f} = D \frac{dC(z)}{dz} - J_{w} C(z) \]  

(7)

At steady state, \( J_{i}^{s} = J_{i}^{f} \)

Or

\[ B(C_{D,m} - C_{F,m}) = D \frac{dC(z)}{dz} - J_{w} C(z) \]  

(8)

With boundary conditions:

\[ \begin{align*}
  z &= 0, C(z) = C_{F,m} \\
  z &= -\delta_{F}, C(z) = C_{F,b}
\end{align*} \]

Integrating (8) with respect to the boundary conditions above, yields

\[ C_{F,m} = C_{F,b} \exp \left( \frac{J_{w}}{k_{F}} \right) + \frac{B \Delta C_{m}}{J_{w}} \left[ \exp \left( \frac{J_{w}}{k_{F}} \right) - 1 \right] \]  

(9)

Subtracting eqn. (9) to (7), \( \Delta C_{m} \) can be derived as:

\[ \Delta C_{m} = (C_{D,m} - C_{F,m}) = C_{D,b} \exp \left[ -J_{w} \left( \frac{1}{k_{D}} + \frac{S}{D_{D}} \right) \right] - C_{F,b} \exp \left( \frac{J_{w}}{k_{F}} \right) + \frac{B \Delta C_{m}}{J_{w}} \left[ \exp \left[ -J_{w} \left( \frac{1}{k_{D}} + \frac{S}{D_{D}} \right) \right] - \exp \left( \frac{J_{w}}{k_{F}} \right) \right] \]  

(10)

or

\[ \Delta C_{m} = \frac{C_{D,b} \exp \left[ -J_{w} \left( \frac{1}{k_{D}} + \frac{S}{D_{D}} \right) \right] - C_{F,b} \exp \left( \frac{J_{w}}{k_{F}} \right)}{1 + \frac{B}{J_{w}} \left\{ \exp \left( \frac{J_{w}}{k_{F}} \right) - \exp \left[ -J_{w} \left( \frac{1}{k_{D}} + \frac{S}{D_{D}} \right) \right] \right\}} \]  

(11)

Assuming van’t Hoff theory is valid in this concentration zone where osmotic pressure is linearly proportional to the salt concentration, the final equation for calculating osmotic water flux in the FO mode becomes:

\[ J_{w}^{FO} = A \frac{\pi_{D,b} \exp \left[ -J_{w} \left( \frac{1}{k_{D}} + \frac{S}{D_{D}} \right) \right] - \pi_{F,b} \exp \left( \frac{J_{w}}{k_{F}} \right)}{1 + \frac{B}{J_{w}} \left\{ \exp \left( \frac{J_{w}}{k_{F}} \right) - \exp \left[ -J_{w} \left( \frac{1}{k_{D}} + \frac{S}{D_{D}} \right) \right] \right\}} \]  

(12)
2.1.2. For the PRO Mode

Likewise, following a similar derivation the complete water flux equation for the PRO mode, including the non-zero hydraulic pressure gradient across the selective layer, can be achieved as follows,

\[
J_{w,PRO} = A \left[ \frac{\pi_{D,b} \exp \left( -\frac{J_w}{k_D} \right) - \pi_{F,b} \exp \left( J_w \left( \frac{1}{k_F} + \frac{S}{D_F} \right) \right)}{1 + \frac{B}{J_w} \exp \left( J_w \left( \frac{1}{k_F} + \frac{S}{D_F} \right) - \exp \left( -\frac{J_w}{k_D} \right) \right)} \right] - \Delta P \quad (13)
\]

Where \( \Delta P \) is the hydraulic pressure applied on the draw side to retard the osmotic water flux from the feed, in conjunction with a hydroturbine providing the means to extract mechanical work from osmosis [39]. The complete derivation is provided in the supplementary material available online.

2.2. New derivation of solute resistivity in FO and PRO

From eqn. (12) and (13), the solute resistivity, \( K_D \) and \( K_F \) in FO and PRO, respectively, can be derived as:

In FO,

\[
K_D = \frac{S}{D_D} = -\frac{1}{k_D} - \frac{1}{J_w} \ln \left[ \frac{J_w}{A} + \frac{B}{A} + \frac{\pi_{D,b}}{k_D} \right] \quad (14)
\]

In PRO,

\[
K_F = \frac{S}{D_F} = -\frac{1}{k_F} + \frac{1}{J_w} \ln \left[ \left( \frac{B}{A} + \frac{\pi_{D,b}}{k_F} \right) \exp \left( -\frac{J_w}{k_D} \right) - \Delta P \left( \frac{1}{J_w} \exp \left( J_w \left( \frac{1}{k_F} + \frac{S}{D_F} \right) - \exp \left( -\frac{J_w}{k_D} \right) \right) \right) \right] \quad (15)
\]

In non-pressurized PRO mode, with \( \Delta P = 0 \):
Eqns. (14)-(16) show that the solute resistivity is, in fact, partly dependent on the mass transfer coefficient of the solute presented in the bulk draw and feed solutions. In other words, the mass transfer of the solute to and from the membrane surface does impact on the behavior of solute molecules inside the porous support layer. This means the increased mass transfer at the porous surface will lead to increased mass transfer within the porous support in both FO and PRO modes. This is well matched with Beavers’s study [53] showing that the increase in mass flow rate through the channel is accompanied by an increase in mass flow rate through the permeate wall (e.g. porous membrane).

2.3. Determination of reflection coefficient and the total resistance of membrane

Reflection coefficient $\sigma$ ($0 < \sigma < 1$) is a measure of the solute reflectivity [54]. $\sigma$ is equal to 0 when the solute and solvent cross the membrane in the same concentration ratio as they are in the adjacent bulk solution [54-56] and is equal to 1 when the solute is completely rejected. In general, $\sigma$ is assumed to be 1 in most studies on FO. However, in reality, no membrane is perfectly selective. Therefore, being able to quantify membrane reflection coefficient could lead to better predictive models of osmotic flux performance and better designs of FO membranes.

In eqn. (12), the osmotic water flux in FO is a function of the water permeability coefficient $A$, the osmotic pressure difference across the membrane $\Delta \pi_m$, and the reverse solute permeability term in the denominator. The reflection coefficient in FO for membranes that are not perfectly selective can therefore be defined as (see Appendix 1 for the derivation):

$$\sigma = \frac{1}{1 + \frac{B}{J_w} \left\{ \exp \left( \frac{J_w}{k_F} \right) - \exp \left[ -J_W \left( \frac{1}{k_D} + \frac{S}{D_D} \right) \right] \right\}}$$

(17)
By expansion of exponential terms in eqn. (17), we then obtain:

\[
\sigma \approx \frac{1}{1 + B \left( \frac{1}{k_F} + \frac{1}{k_D} + K \right)} = \frac{1}{B} \frac{1}{1 + B \left( \frac{1}{k_F} + \frac{1}{k_D} + K \right)}
\]

Likewise, the reflection coefficient in PRO can be derived from eqn. (13):

\[
\sigma = \frac{1}{1 + \frac{B}{J_w} \left\{ \exp \left[ J_w \left( \frac{1}{k_F} + \frac{S}{D_F} \right) - \exp \left( -\frac{J_w}{k_D} \right) \right] \right\}}
\]

After expansion, we have:

\[
\sigma \approx \frac{1}{1 + B \left( \frac{1}{k_F} + \frac{1}{k_D} + K \right)} = \frac{1}{B} \frac{1}{1 + B \left( \frac{1}{k_F} + \frac{1}{k_D} + K \right)}
\]

Appendix A2 shows the error of the approximation by Taylor expansion. In eqns. (18) and (20), the denominator is the sum of the solute resistivity due to (1) membrane selective layer, (2) support structure and (3) mass transfer boundary layers at membranes surfaces. If one analogizes the diffusive flow rate of the solute across the membrane to an electric current through a passive resistor, the total resistance of membrane to solute transport could be defined as the sum of the intrinsic membrane resistance and the solution resistance affected by the convective rate in the bulk fluid [63]. Equations (18) and (20) have similarity to the total solute resistance across the membrane which is defined as,

\[
R = \frac{1}{B} \frac{1}{k_F} + \frac{1}{k_D} + K
\]
we can then see that, in osmosis, the membrane reflection coefficient is the percentage of the solute resistivity, $1/B$, of the membrane selective layer in the total resistance of the membrane, $R$. Reflection coefficient is dependent on both membrane characteristics (selectivity, structural parameter) and solution properties (diffusivity, hydrodynamic conditions). This is in good agreement with Anderson and Malone’s study showing that reflection coefficient is a function of solute and membrane characteristics [55]. Note that in equation (21), we assumed an insignificant resistance arisen at the membrane-solution interface.

3. Materials and Methods

3.1. Materials

Sodium chloride was purchased from Fisher Scientific (Pittsburgh, PA). Isopropanol was purchased from J.T. Baker (Center Valley, PA). Water used in this study was ultrapure Milli-Q water produce by a Millipore Integral 10 water system, (Millipore Corporation, Billerica, MA).

3.2. Membranes

Two different types of asymmetric membranes specifically designed for osmotically driven membrane processes were used in this study. These were an asymmetric cellulose triacetate membrane embedded with a nonwoven mesh from Hydration Technology Innovations (HTI, Albany, OR) and an early generation of thin-film composite (TFC) membrane from Oasys Water (March 2011, Boston, MA). These two membranes are denoted as HTI and Oasys, respectively. The structure of HTI’s cellulose triacetate membrane has been broadly studied and reported elsewhere [38, 57]. Oasys TFC membrane structure and chemistry are proprietary and have not been widely explored. It is believed to have a thin polyamide selective layer deposited on a porous support which had been casted on a nonwoven backing layer [64].
3.3. Membrane performance characteristics

3.3.1. Intrinsic membrane properties determined by reverse osmosis

The water permeance of these membranes was measured in a lab scale reverse osmosis testing system at pressures of 10.3, 15.5, 20.7, 25.9 and 31.0 bar (150, 225, 300, 375 and 450 psi) at temperatures of 20°C and 40°C and ambient RH. Fresh membrane samples were used to determine water permeability coefficients at each of the tested temperatures. Rejection tests were carried out following the measurement of water permeability coefficients at 15.5 bar with a 2000 ppm NaCl feed at 20°C and 40°C using a cross flow velocity of 0.25 m/s. NaCl rejection was measured using a conductivity. Based upon hydrodynamic conditions of the system and empirical data from Comesana [58] and Lobo [59], intrinsic rejection was determined from a Sherwood number correlation [1]. Intrinsic rejection was used to determine the sodium chloride permeability for this membrane, calculated from the following equation [1].

\[
B = \frac{(1-R)A(\Delta P - \Delta \pi)}{R} = \frac{(1-R)J_w}{R}
\]

Here B is the solute permeability coefficient, R is the rejection, A is the water permeability coefficient of the membrane, \(\Delta P\) is the transmembrane hydrostatic pressure, \(\Delta \pi\) is the transmembrane osmotic pressure, and \(J_w\) is the water flux of the rejection measurement.

3.2.2 FO/PRO osmotic flux performance

Cross-flow FO and PRO tests were carried out at 20 °C and 40 °C in a lab-scale FO system described elsewhere [25]. No spacers were used. Sodium chloride (NaCl) was used as a model draw and feed solute for this study. Osmotic water flux was measured in both FO (the porous support faces the draw) and PRO (the porous support faces the feed) modes. Zero transmembrane pressure was maintained. For FO, deionized (DI) water was used as the feed solution, whereas 0.5 M, 1.0 M and 1.5 M NaCl were used as the draw solutions. For PRO, 0.05
M, 0.5 M and 1.0 M NaCl solutions were used in the feed, while 1.5 M NaCl solution was fixed as the draw. The feed and draw solutions were maintained the same temperature, pressure and cross-flow velocity. Solution properties (e.g. density, viscosity, diffusivity, osmotic coefficient) varying at different concentrations and temperatures were obtained from literature [58, 59] and presented in Table 1. Membranes were stored in DI water at 4 °C prior to testing without pretreatment. Cross-flow velocities used for testing ranged from 10.3 to 20.7, 31.0, 41.4 and 46.5 cm/s and were maintained for both solutions. The experimental protocol was described in detail in section 2.4 of McCutcheon et al. [38].

4. Results and Discussion
4.1. Membrane permselectivity and osmotic flux performance

Intrinsic water and solute permeance A and B of HTI and Oasys membranes from RO tests are tabulated in Table 2. In general, A and B increased with feed temperature. These values were then substituted into the flux model, eqns. 12 and 13. Figure 2 illustrates the iteration procedure using a Matlab algorithm.

Figure 3 shows an increase of osmotic water flux with cross-flow velocity in both FO and PRO experiments. The primary reason for this increased flux is increased mass transfer coefficient and a reduced ECP. Dilution factor of the draw solution and concentration factor of the feed could be considered as well (since crossflow velocity impacts residence time), but these effects were negligible in our system because of the small membrane area and very short residence time (less than 1 second in the cell). In a module, there might be different effects, but our tests were limited to coupons. Figure 3 shows that, both membranes demonstrated significantly enhanced water flux when cross-flow velocities increased from 10.3 to 20.7 cm/s.
The data shows that the HTI cellulosic membrane has substantially lower flux than the Oasys membrane for both FO and PRO mode. The Oasys membrane performed 2 – 3 times higher fluxes at the lowest cross-flow velocity than the HTI membrane (Figure 3). This means that all polarization effects will be less severe for the HTI membrane. Figure 3a shows an approximate 40% increase in water flux over the range of crossflow velocities tested for HTI membrane in both the FO and PRO modes. However, the fluxes for the Oasys membrane were nearly doubled over this same range. Flux is more sensitive to crossflow velocity when the concentration polarization is more severe.

4.2. Quantifying Concentration Polarization

ECP modulus is defined as a ratio of the osmotic pressure measured at the membrane surface to that obtained in the bulk solution [38]. ICP modulus, on the other hand, is the ratio of the osmotic pressure at the support – selective layer interface to that measured at the membrane surface [38]. Dilutive ECP/ICP modulus is between 0 and 1 whereas concentrative ECP/ICP modulus is greater than 1 [38]. The farther these ratios are from the value of 1 is indicative of more severe polarization.

4.2.1. External concentration polarization on the porous side (ECP\textsuperscript{p}) of EO membranes

Figure 4 shows the dilutive and concentrative ECP\textsuperscript{p} moduli [38] obtained from FO and PRO tests at different draw and feed concentrations, respectively, for HTI and Oasys membranes. ECP moduli were calculated from the exponential function of ratio of water flux to mass transfer coefficient [38, 42]. Results show that these moduli deviate from 1 substantially in both cases, and therefore, should not be neglected in the prediction of FO membrane performance (as they are currently). The Oasys TFC membrane shows a more severe ECP\textsuperscript{p}
compared to the HTI membrane due to the fact that the membrane performs at higher water flux at the same conditions.

Water flux was found to be the key factor affecting the severity of ECP\textsuperscript{p}. In the FO mode, ECP\textsuperscript{p} became more severe when the draw concentration was increased from 0.5 M to 1.0 M and 1.5 M NaCl (Figure 4). However, the decrease of concentrative ECP\textsuperscript{p} with increasing feed concentration from 0.05 M to 0.5 M and 1.0 M NaCl, in PRO, was more noticeable, especially for the Oasys membrane. Increasing feed concentration significantly reduced osmotic driving force leading to a reduction in water flux and a decrease in ECP\textsuperscript{p}.

4.2.2. Internal concentration polarization (ICP)

4.2.2.1. Impacts of Crossflow Velocity on ICP

Figure 5 shows the dilutive ICP modulus in the FO mode for both membranes calculated based on eqn. (15) listed elsewhere [38]. As noted by the increasing modulus, ICP became less severe with increasing cross-flow velocity. This is counter to the commonly belief that ICP is an unstirred boundary layer inside the porous support and, therefore, cannot be mitigated by hydrodynamic conditions [1, 39, 60]. In fact, crossflow velocity induces mixing in the porous support layer, especially for well-designed membranes that have thin and highly porous support layers. Such an effect was previously observed by Beavers and Joseph, where an increase in mass flow rate through the channel is accompanied by an increase in mass flow rate through a porous wall (e.g. porous membrane) [53]. As shown in Figure 5b, shear effects are transmitted into the body of the porous materials through a boundary layer region [53]. It indicates that increasing the cross-flow velocity not only thins the ECP\textsuperscript{p} boundary layer, it changes the slip velocity of fluid at membrane support layer surface. This impacts mass transfer into the support layer itself, especially at higher crossflow velocities. Other properties, such as high surface porosity or roughness may intensify this effect. Note that both HTI and Oasys membrane
supports may have relatively high surface roughness and high porosity due to the large average pore diameter of the support layer (i.e. the embedded mesh and polyester non-woven fabric for HTI and Oasys, respectively). Results show that the HTI and Oasys membranes behaved differently with increased cross-flow velocity. For the Oasys membrane, dilutive ICP modulus increased substantially when velocity increased from 10.3 to 20.7 cm/s and then plateaued. On the other hand, for the HTI membrane, ICP modulus increased gradually and continuously across the full range of cross-flow velocities considered. These results indicate that, for the more optimized support structure of the Oasys membrane, ICP can be diminished at high crossflow velocity, though this benefit has diminishing returns at even higher velocity. This is not the case for the asymmetric HTI membrane structure which consists of a polyester mesh embedded in a porous polymer layer. This denser structure resists internal mixing and thus the benefits are improved throughout a wide range of crossflow velocities.

4.2.2.2. Impact of draw and feed concentrations on ICP modulus

The impact of draw and feed concentration in FO and PRO has long been studied across the osmotic separations literature. However, little has been done with examining the impact of these solution concentrations with relation to concentration polarization both inside and outside of the membrane structure. Our results (Figure 6) definitively show that ICP modulus is more severe with increased draw concentration in FO and decreased feed concentration in PRO. Also, ICP impacted the Oasys membrane more intensely than the HTI membrane, especially in FO mode. This is largely due to the substantially higher fluxes generated with the Oasys Membrane.
4.2.2.3. Comparison with previous flux models

Figure 6 also compares our work to earlier McCutcheon’s and Yip-Tiraferti’s flux models. This work shows a less severe ICP since it no longer incorporates ECP along the support layer side of the membrane (ECP<sub>p</sub>). Our data shows that these earlier estimates led to an overprediction of the ICP modulus by 20 – 50%. The overprediction is more pronounced for higher flux membranes, such as the Oasys membrane. Among these data, one of the largest discrepancies (~50%) occurred for the Oasys membrane at the lowest feed concentration of 0.05M.

4.3. Lost driving force caused by concentration polarization

While the modulus is an important value when describing boundary layer phenomenon with FO, it may not be representative of the true impact of the polarization on membrane performance. For example, a high concentrative modulus may have little impact on the driving force if the solute concentration is very low. We therefore must consider the real impact on driving force caused by concentration polarization. Figures 7 and 8 show the lost driving force caused by concentration polarizations in FO and PRO modes, respectively, as a percentage of total lost driving force. In FO, although ICP is responsible for most of the loss of driving force, ECP<sub>p</sub> is responsible for ~20% and ~40% of the driving force loss for the HTI and Oasys membranes, respectively. Since DI water was used as the feed solution, ECP on the feed was negligible and therefore not visible in the graph (Figure 7). In the PRO mode, the ICP contributed less to the reduction in driving force. Instead, ECP<sub>s</sub> (on the selective side) significantly reduced the driving force (up to 40% for the Oasys membrane). Although ECP<sub>p</sub> was less pronounced, it was non-negligible (~8% for the HTI and ~20% for the Oasys membrane). In
all cases, the Oasys TFC membrane support appeared to be significantly impacted by the ECP boundary layers at membrane surfaces. This is expected given the high flux of these membranes. The HTI membrane performance was severely influenced by the ICP within the membrane support structure.

4.4. Solute resistivity, $K$

Figure 9 shows that solute resistivity, $K$, obtained from FO was higher than PRO due to the more severe ICP within the support. As commonly defined, $K$ is an intrinsic property of membrane support for a given draw solution (at a specific diffusivity). It is, therefore, believed to be unaffected by hydrodynamic conditions at the surface of the support layer. However, in Figure 9, $K$ is shown to decrease with increasing cross-flow velocity in both membranes and orientations. These results are in agreement with equations (14) and (16) showing that $K$ is a function of mass transfer coefficient, $k$, at the surface of the support. As mentioned, tangential cross-flow velocity can partially induce mixing within the porous support leading to increased effective diffusivity of solutes (Reynolds numbers and mass transfer coefficients in the draw and the feed are tabulated in Table 3). As such, the porous support is no longer an unstirred zone and the original assumptions used in defining solute resistivity are no longer valid [39].

4.5. The true structural parameter

In general, empirical structural parameter measured in the PRO mode was smaller than from FO (Figure 10). Results show that structural parameters obtained from McCutcheon’s and Yip-Tiraferri’s flux models were higher than those calculated from this study by 15 – 45 %. Note that unlike many previous models, we took into account the difference of mass transfer coefficients between the feed and the draw sides, e.g. $k_F$ and $k_D$, respectively. Also,
dependences of the solution physical properties (e.g. density, viscosity, diffusivity, etc.) on the testing conditions (e.g. concentration, temperature) were included in our study.

Furthermore, calculated structural parameters decreased with cross-flow velocity. Originally, we had hoped that proper accounting of the mass transfer resistance would have resulted in the S parameter being calculated as being the same regardless of the testing conditions. However, due to the mixing induced in the support layer caused by increased tangential flow as suggested in Figure 5b, the underlying assumption of a stagnant layer in the support layer is incorrect. This means that the “no slip” condition (where crossflow is zero) is actually within the support layer and this reduces the thickness of the support layer, perhaps by a substantial amount. Thus, since the model assumes “no mixing” in the support layer, and mixing that does occur will force the model to output lower structural parameters to compensate.

Likewise, the structural parameter has been theoretically derived as an effective thickness of the membrane support and should therefore be independent of operating conditions such as membrane orientation, solution concentration, and temperature. Park et al. [48] presented a numerical model based on finite element method to determine a relatively constant membrane structural parameter in FO. However, as empirical S values have been broadly estimated and reported using common flux models [21, 38, 52], understanding on how and why the value of this parameter changes with testing conditions is imperative. This is even more important considering few groups around the world use a standardized method for measuring structural parameter [61].

In fact, changes in empirical S with testing conditions partly dictate transport behaviors of solute within the membrane structure. Figure 11 shows structural parameters of the HTI and Oasys membranes calculated across different operating temperatures and draw and feed concentrations in FO and PRO modes using our model. In general, we saw relatively constant
structural parameter for the Oasys TFC membrane over a broad range of solutions conditions. However, structural parameter values for the HTI membranes changed substantially with varying testing conditions. The HTI membrane S values changed with increased draw concentration in FO while they gradually decreased with feed concentration in PRO.

4.6. Total resistance of membranes to solute transport, R

As defined in eqn. (21), total resistance to solute transport (R) is the sum of all resistivity that solute encounters when transporting across the membrane from the bulk draw to the bulk feed solutions. These include solute resistivity due to membrane selective layer (1/B), the support layer (K) and the mass transfer boundary layers at membrane surfaces (1/k_D and 1/k_F).

R can be considered as a combination of membrane resistance and solution resistance to ions transport (see Figure 1). Total resistance is impacted by membrane structure and testing conditions (e.g. cross-flow velocity, temperature). It is, therefore, necessary to improve membrane design to obtain a minimal total resistance to solute transport toward the selective layer while also adjusting our testing and data analysis protocols to ensure that we are accurately attributing performance metrics to the appropriate membrane property.

Figure 12 shows that total solute resistance of the Oasys TFC membrane is less than that of HTI asymmetric membrane due mainly to its smaller solute resistivity K. In addition, increasing cross-flow velocity reduced total membrane solute resistance because it increased mass transfer at membrane surfaces, thinned the ECP boundary layers, and reduced ICP. However, the insignificant decrease of total resistance with further increasing cross-flow velocities indicates that the membrane intrinsic parameters affect solute resistance more than process hydrodynamics. Furthermore, R was found to decrease substantially with increasing temperature. Raising testing temperature from 20 to 40°C led to increases in solute diffusivity
(by ~35%, table 1), mass transfer coefficients \( k_D \) and \( k_F \) as a function of crossflow velocity (by ~12 – 40%), and the solute permeability coefficient \( B \) (by ~6% for the HTI membrane, ~30% for the Oasys membrane, table 2). Therefore, the decrease of \( R \) with temperature was more pronounced for the Oasys membrane (~33%) than the HTI membrane (~10%).

### 4.7. Reflection coefficient

In eqns. (18) or (20), we defined reflection coefficient, \( \sigma \ (0 < \sigma < 1) \), as the ratio of the solute resistivity of the membrane selective layer and the total resistance to solute transport, \( R \). It is inversely proportional to total resistance and is, therefore, determined by both membrane characteristics (selectivity, structural parameter) and solute behaviors (diffusivity, mass transfer coefficient). Table 4 tabulates reflection coefficient of Oasys and HTI membranes at different testing conditions (e.g. membrane orientation, draw and feed concentrations, and cross-flow velocity). In general, the HTI membrane, with its smaller \( B \), was more selective than the Oasys TFC membrane.

It is hypothesized that solute behaviors are dependent on the coupling effects of (1) the tangential fluid flows at membrane surfaces, (2) the convective water flow and (3) the diffusive solute flow across the membrane. A change in intensity of any flow would impact the *residence time* of solutes in the vicinity of membrane surfaces. The longer the solutes stay at membrane surfaces, the more likely these molecules would partition to the selective layers and diffuse across the membranes. As such, membranes become more selective at shorter solute *residence time*. In FO, the increased cross-flow velocity along with increased mixing and enhanced convective water flux would swiftly sweep the solutes away from membrane surfaces. Due to the reduced *residence time*, solutes had less chance to partition and penetrate across the membrane. Therefore, membranes tend to be more selective at high water flux (i.e. high cross-flow velocity and high draw concentration).
On the other hand, in PRO, the convective water flux, while carrying solute molecules from the bulk feed, is driven toward the support – selective layer interface. As a result, lower water flux would reduce the chances of solutes to stay in the vicinity of selective layer surface to diffuse through. As the feed concentration increased from 0.05 M to 1 M NaCl, the osmotic water flux reduced significantly leading to an increased reflection coefficient. Namely, the membrane was more selective with increased feed concentration. However, in PRO, when mixing is introduced at high cross-flow velocity, there is a tradeoff between mixing intensity and the convective water flux. Increased cross-flow velocity led to enhanced osmotic water flux. However, this enhancement of water flux was not sufficient to surpass the mixing intensity generated by the tangential fluid flow. It resulted in a shorter residence time for solutes to stay within the porous support and thus increased reflection coefficient at high cross-flow velocity.

5. Conclusions

As noted throughout this paper, there are numerous complexities associated with understanding mass transfer during osmotic flux across asymmetric membranes. While internal concentration polarization has been studied in detail for some time, rarely has it been placed in context with all other resistances for solute flux (which ultimately dictates driving force in osmosis). The model introduced here allows for a full accounting, and deconvolution, of all resistances present for solute transport during osmosis. This allows for a more accurate calculation of typical membrane design metrics, such as structural parameter. Accuracy in these calculations is critical if membrane and system designers are expected to increase process efficiency. If, for instance as this study shows, structural parameter is overpredicted, then designers will wrongly focus on membrane design, only to find out later that hydrodynamics is more important than previously thought. This overprediction becomes more of an issue as better membranes are made (as has been seen recently by numerous academic groups and companies around the world). Accurate accounting of resistances through well-developed
models that are verified by experiments and analysis of membrane structure will be essential as forward osmosis and its sibling processes gain their first footholds into the commercial sector and thereby require extensive optimization.

Acknowledgements

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Nomenclature

\[ \begin{align*}
\text{ECP}^s & : \text{external concentration polarization, at the selective layer surface} \\
\text{ECP}^p & : \text{external concentration polarization, at the porous support} \\
c\text{ECP}, d\text{ECP} & : \text{concentrative/dilutive external concentration polarization} \\
J_w & : \text{osmotic water flux (gal.ft}^2\text{day}^{-1}) \\
J_{s,\alpha} & : \text{reverse salt flux across the selective layer (gal.ft}^2\text{day}^{-1}) \\
J_{s,d}, J_{s,f} & : \text{reverse salt flux across the porous support, facing against the draw/feed solution, respectively (gal.ft}^2\text{day}^{-1}) \\
A & : \text{water permeability coefficient (gal.ft}^2\text{day}^{-1}\text{psi}^{-1}) \\
B & : \text{solute permeability coefficient (gal.ft}^2\text{day}^{-1}) \\
\sigma & : \text{reflection coefficient}
\end{align*} \]
\( \Delta \pi_m \) \hspace{1cm} \text{osmotic pressure difference across the selective layer (bar)}

\( \pi_{D,m}, \pi_{F,m} \) \hspace{1cm} \text{osmotic pressure of the draw/feed solution, at two sides of selective layer (bar)}

\( \pi_{D,b}, \pi_{F,b} \) \hspace{1cm} \text{osmotic pressure of the draw/feed solution, in the bulk (bar)}

\( \pi_{D,i}, \pi_{F,i} \) \hspace{1cm} \text{osmotic pressure of the draw/feed solution, at the surface of the porous support (bar)}

\( \Delta P \) \hspace{1cm} \text{transmembrane hydraulic pressure difference (bar)}

\( \Delta C_m \) \hspace{1cm} \text{concentration gradient across the selective layer (mol/l)}

\( C_{D,m}, C_{F,m} \) \hspace{1cm} \text{concentration of the draw/feed solution, at two sides of selective layer (mol/l)}

\( C_{D,b}, C_{F,b} \) \hspace{1cm} \text{concentration of the draw/feed solution, in the bulk (mol/l)}

\( C_{D,i}, C_{F,i} \) \hspace{1cm} \text{concentration of the draw/feed solution, at the surface of the porous support (mol/l)}

\( D \) \hspace{1cm} \text{solute diffusion coefficient (m}^2/\text{s)}

\( D_s \) \hspace{1cm} \text{solute diffusion coefficient inside the support layer (m}^2/\text{s)}

\( D_D, D_F \) \hspace{1cm} \text{solute diffusion coefficient in the bulk draw/feed (m}^2/\text{s)}

\( x, z \) \hspace{1cm} \text{coordinates perpendicular to membrane axis}

\( C(x), C(z) \) \hspace{1cm} \text{solute concentration varies with positions (mol/l)}

\( \varepsilon \) \hspace{1cm} \text{porosity of membrane support}

\( \tau \) \hspace{1cm} \text{tortuosity of membrane support}

\( t_s \) \hspace{1cm} \text{thickness of the membrane support (µm)}
\( \delta_D, \delta_F \)

thicknes of the boundary layer at membrane surfaces, facing against the draw/feed (\( \mu \text{m} \))

\( S \)

structural parameter of the support layer (\( \mu \text{m} \))

\( k_D, k_F \)

mass transfer coefficient of the draw/feed (\( \text{m/s} \))

\( K_D, K_F \)

solute resistivity of membrane support, facing against draw/feed (\( \text{s/m} \))

\( R \)

membrane total resistance to solute transport (\( \text{s/m} \))
References


Highlights

- A mathematical flux model for osmotic transport was proposed and validated empirically.
- All significant boundary layers were encapsulated and effectively deconvoluted.
- The ECP on the porous side of the membrane cannot be ignored.
- Equations of total solute resistance and reflection coefficient were introduced.
- Membrane selectivity is determined by residence time of solute at membrane surface.
Graphical Abstract
Figure 1 – Schematics of solute concentration profile at steady state across a thin-film composite membrane in (a) FO mode and (b) PRO mode.
Figure 2 – Iteration procedure using Matlab algorithm to de-convolute ECP, ICP moduli, structural parameter, and mass transfer resistance, given operating conditions (i.e. solutions concentration, temperature, cross-flow velocities, channel dimensions and membranes permselectivity obtained from reverse osmosis tests).
Figure 3 – Osmotic water flux performance of HTI and Oasys membranes as a function of cross-flow velocity (in FO and PRO modes) at 20 °C. Similar flux behavior was observed at 40 °C (supplementary information).
Figure 4 – ECP moduli on the porous side of the membranes in FO and PRO modes obtained from this study. The dash lines represent for the ideal case with no concentration polarization. Data was obtained at 20°C and a cross-flow velocity of ~31.0 cm/s.
Figure 5 – (a) Change of ICP with cross-flow velocities representatively in FO process and (b) Velocity profile for the rectilinear flow in a horizontal channel formed by a permeable lower wall (y = 0) and an impermeable upper wall (y = h) [53] (Reprinted by permission of Cambridge University Press).
Figure 6 – Comparing ICP modulus from FO and PRO experiments on HTI asymmetric and Oasys TFC membranes using new model and previous models by McCutcheon [38], Yip [21] and Tiraferri [52]. Data was obtained at 20°C and a cross-flow velocity of ~31.0 cm/s.
Figure 7 - Lost driving force caused by concentration polarizations in FO mode at 20 and 40 °C.

Zero transmembrane pressure difference. Bars represent for percentages of $\frac{\Delta \pi}{\Delta \pi_{\text{theo}}}$. Blue bars represent for the effective driving force across the selective layer. Note that the contribution of the ECP on the selective layer (feed) side (yellow), although non-zero, was insignificant to be visible in the plots (ranging from 0.2 – 0.65%, see Appendix 3). Please refer to online version to view this image in color.
Figure 8 - Lost driving force caused by concentration polarizations in PRO mode at 20 and 40 °C. Zero transmembrane pressure difference. Bars represent for percentages of \( \frac{\Delta \pi_i}{\Delta \pi_{\text{theo}}} \) (i = 1 – 4). Blue bars represent for the effective driving force across the selective layer.

Please refer to online version to view this image in color.
Figure 9 – Solute resistivity, K, of HTI asymmetric and Oasys TFC membranes at different crossflow velocities and temperatures. Data was averaged from independent tests of different membrane coupons across the ranges of solution concentration.
Figure 10 – Structural parameter of HTI and Oasys TFC membranes vs. cross-flow velocity in FO and PRO tests obtained from different models. Data was averaged from independent tests of different membrane coupons.
Figure 11 - Impacts of draw and feed solutions temperature and concentrations on the structural parameter of HTI and Oasys TFC membranes in both FO and PRO modes.

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Figure 12 – Changes of total resistance of membranes to solute transport with cross-flow velocity and testing temperature. Eqn. (21) was used to obtain total resistance to solute transport. Values were averaged from independent data of different membrane coupons across the ranges of solution concentration in both FO and PRO modes.
Table 1 – Physical properties of fluids and mass transfer coefficient $k$ of the draw and the feed at different temperature and concentrations [58, 59].

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<th>Concentration (M)</th>
<th>Density $(\text{kg/m}^3)$</th>
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Table 2 – Membranes permselectivity tested in RO at 20 °C and 40 °C. Water permeability coefficient, $A$, was obtained at pressures from 150 to 450 psi. Solute permeability coefficient, $B$ was determined at 225 psi, using 2000 ppm NaCl feed solution.

<table>
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<tr>
<th>Temperature (°C)</th>
<th>HTI ($A$ (LMH/bar))</th>
<th>HTI ($B$ (LMH))</th>
<th>Oasys ($A$ (LMH/bar))</th>
<th>Oasys ($B$ (LMH))</th>
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Table 3 – Reynolds numbers and predicted mass transfer coefficients $k_D$, $k_F$, at different testing conditions

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Table 4 - Reflection coefficient of membranes in FO and PRO tests.

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<th>Crossflow Velocity (cm/s)</th>
<th>FO - DI water feed</th>
<th>PRO - 1.5M NaCl draw</th>
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